

A NONSTANDARD SOLUTION TO THE WAVE EQUATION

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ABSTRACT. We follow d'Alembert's method for solving the standard wave equation, in the nonstandard case, and compare the solutions.

Definition 0.1. Given $u \in V(\overline{\mathcal{R}}_\eta \times \overline{\mathcal{T}}_\nu)$, we let;

$$u_t(y, s) = \nu(u(y, s + \frac{1}{\nu}) - u(y, s)), \text{ for } (y, s) \in \overline{\mathcal{R}}_\eta \times (\overline{\mathcal{T}}_\nu \setminus \frac{\nu^2-1}{\nu})$$

$$u_x(y, s) = \frac{\sqrt{\eta}}{2}(u(y + \frac{1}{\sqrt{\eta}}, s) - u(y - \frac{1}{\sqrt{\eta}}, s)), \text{ for } (y, s) \in \overline{\mathcal{R}}_\eta \times \overline{\mathcal{T}}_\nu$$

We define the nonstandard unconvoluted wave equation, on $\overline{\mathcal{R}}_\eta \times (\overline{\mathcal{T}}_\nu \setminus \frac{\nu^2-1}{\nu}, \frac{\nu^2-2}{\nu})$ by;

$$u_{tt} - u_{xx} = 0$$

Lemma 0.2. Given initial conditions $\{f, g\} \subset V(\overline{\mathcal{R}}_\eta)$ there exists a unique $u \in V(\overline{\mathcal{R}}_\eta \times \overline{\mathcal{T}}_\nu)$ solving the nonstandard unconvoluted wave equation, such that $u^0 = f$, and $u_t^0 = g$, and;

$$u(x, t + \frac{2}{\nu}) = 2u(x, t + \frac{1}{\nu}) + \frac{\eta}{4\nu^2}u(x + \frac{2}{\sqrt{\eta}}, t) - (1 + \frac{\eta}{2\nu^2})u(x, t)$$

$$+ \frac{\eta}{4\nu^2}u(x - \frac{2}{\sqrt{\eta}}, t), (*)$$

$$\text{for } (x, t) \in \overline{\mathcal{R}}_\eta \times (\overline{\mathcal{T}}_\nu \setminus \frac{\nu^2-1}{\nu}, \frac{\nu^2-2}{\nu})$$

Moreover, if the standard initial conditions $\{u^0, u_t^0\}$ are given with $\frac{\partial^i(u^0)}{\partial x^i} \leq D$ and $\frac{\partial^j(u_t^0)}{\partial x^j} \leq D$, uniformly, for finite $(i, j) \in \mathcal{N}^2$, then for $(t_0, x_0) \in \overline{\mathcal{T}}_\nu \times \overline{\mathcal{R}}_\eta$, with t_0 finite, and the choice $\eta \leq 4\nu$, $u_{txx}|_{t_0, x_0}$ and $u_{xx}|_{t_0, x_0}$ are bounded, for the nonstandard equation generated by initial conditions $\{v, v_t\}$, with $v^0 = (u^0)_\eta$, $v_t^0 = (u_t^0)_\eta$.

Proof. Using the definition of the derivatives, the equation in Definition 0.1, and rearranging, we obtain the defining schema for u given in

(*). We are free to choose the values for the first two time steps, by setting;

$$u(x, 0) = f(x)$$

$$u_t(x, 0) = \nu(u(x, \frac{1}{\nu}) - u(x, 0)) = g(x)$$

First, observe, using Taylors' Theorem, that $v_{x^n}^0$ and $v_{tx^n}^0$ are bounded. Suppose inductively, that there exists a constant $C_i \in \mathcal{R}$, for $\frac{i}{\nu}$ finite, with;

$$\max(\{u_{\frac{j}{\nu}, x^n}, u_{\frac{k}{\nu}, x^n} : 0 \leq j \leq i, 0 \leq k \leq i-1\}) \leq C_i$$

for n even, where v_{x^n} denotes the n 'th derivative of v with respect to x .

We have that;

$$\begin{aligned} v(x, \frac{i+1}{\nu}) &= 2v(x, \frac{i}{\nu}) + \frac{\eta}{4\nu^2}v(x + \frac{2}{\sqrt{\eta}}, \frac{i-1}{\nu}) - (1 + \frac{\eta}{2\nu^2})v(x, \frac{i-1}{\nu}) \\ &+ \frac{\eta}{4\nu^2}v(x - \frac{2}{\sqrt{\eta}}, \frac{i-1}{\nu}) \end{aligned}$$

and taking the n 'th even derivative with respect to x ;

$$\begin{aligned} v_{x^n}(x, \frac{i+1}{\nu}) &= 2v_{x^n}(x, \frac{i}{\nu}) + \frac{\eta}{4\nu^2}v_{x^n}(x + \frac{2}{\sqrt{\eta}}, \frac{i-1}{\nu}) - (1 + \frac{\eta}{2\nu^2})v_{x^n}(x, \frac{i-1}{\nu}) \\ &+ \frac{\eta}{4\nu^2}v_{x^n}(x - \frac{2}{\sqrt{\eta}}, \frac{i-1}{\nu}) \end{aligned}$$

Abbreviating notation;

$$v_{i,tx^n} = \nu(v_{i+1,x^n} - v_{i,x^n}) = \nu(2v_{i,x^n} - v_{i-1,x^n} + \frac{\eta}{4\nu^2}(v_{i-1,x^n}^{lsh^2} - 2v_{i-1,x^n} + v_{i-1,x^n}^{rsh^2}) - v_{i,x^n})$$

$$\nu(v_{i,x^n} - v_{i-1,x^n}) + \frac{\eta}{4\nu}(v_{i-1,x^n}^{lsh^2} - 2v_{i-1,x^n} + v_{i-1,x^n}^{rsh^2})$$

$$= v_{i-1,tx^n} + \frac{1}{4\nu}v_{i-1,x^{n+2}}$$

$$|v_{i,tx^n}| \leq |v_{i-1,tx^n}| + \frac{1}{4\nu}|v_{i-1,x^{n+2}}|$$

$$\leq C + \frac{1}{4\nu}C = C(1 + \frac{1}{4\nu})$$

$$\begin{aligned}
v_{i+1,x^n} &= 2v_{i,x^n} + \frac{\eta}{4\nu^2}(v_{i-1,x^n}^{lsh^2} - (1 + \frac{\eta}{2\nu^2})v_{i-1,x^n} + \frac{\eta}{4\nu^2}v_{i-1,x^n}^{rsh^2}) \\
&= v_{i,x^n} - v_{i-1,x^n} + v_{i,x^n} + \frac{\eta}{4\nu^2}(v_{i-1,x^n}^{lsh^2} - 2v_{i-1,x^n} + v_{i-1,x^n}^{rsh^2}) \\
&= \frac{v_{i-1,tx^n}}{\nu} + v_{i,x^n} + \frac{\eta}{4\nu^2}(v_{i-1,x^{n+2}}) \\
|v_{i+1,x^n}| &\leq \frac{C}{\nu} + C(1 + \frac{\eta}{4\nu^2})
\end{aligned}$$

Iterating, we obtain that;

for finite t ;

$$\begin{aligned}
|v_{t,tx^n}| &\leq C(1 + \frac{1}{4\nu})^{[\nu t]} \\
&= C(1 + \frac{1}{4\nu})^{\nu \frac{[\nu t]}{\nu}} \\
&\simeq Ce^{4\frac{[\nu t]}{\nu}} \simeq Ce^{4t} \\
|v_{t,x^n}| &\leq C(1 + \frac{1}{\nu} + \frac{\eta}{4\nu^2})^{[\nu t]} \\
&= C(1 + (\frac{1+\frac{\eta}{4\nu}}{\nu})^{[\nu t]}) \\
&\leq C(1 + (\frac{2}{\nu})^{[\nu t]}) \\
&\simeq Ce^{2t}
\end{aligned}$$

with $\eta \leq 4\nu$, as required. A similar argument works for the odd derivatives $\{u_{x^n}, u_{tx^n}\}$, with $n \in \mathcal{N}$ odd.

□

Lemma 0.3. *Suppose that u satisfies the nonstandard equation in Lemma 0.2, with the extra assumption that for finite $(i, j) \in \mathcal{N}^2$, $\frac{\partial^i(u^0)}{\partial x^i}$ and $\frac{\partial^j(u_t^0)}{\partial x^j}$ are bounded, for finite $(x, t) \in \overline{\mathcal{R}}_\eta \times \overline{\mathcal{T}}_\nu$. Then, for such (x, t) ;*

$$u(x, t) \simeq \frac{1}{2}(u(x+t, 0) + u(x-t, 0)) + \int_{[x-t, x+t]} u_t^0(w) d\mu_\eta(w)$$

In particular, if F satisfies the standard wave equation on $\overline{\mathcal{R}}^2$, with the property that $F^0, F_t^0 \subset C^\infty(\mathcal{R})$, then if u satisfies the nonstandard equation with initial conditions $\{(F^0)_\eta, (F_t^0)_\eta\}$, we have that u is S -continuous for finite $(x, t) \in \overline{\mathcal{R}}_\eta \times \overline{\mathcal{T}}_\nu$, and;

$$\circ u(x, t) = F(\circ x, \circ t)$$

Proof. By the previous lemma, all finite derivatives of the form u_{x^n} and u_{tx^n} are uniformly bounded. Factoring the equation $u_{tt} - u_{xx} = 0$, we obtain that;

$$\left(\frac{\partial}{\partial t} + \frac{\partial}{\partial x}\right)\left(\frac{\partial}{\partial t} - \frac{\partial}{\partial x}\right)u = 0$$

for $(x, t) \in \overline{\mathcal{R}_\eta} \times (\overline{\mathcal{T}_\nu} \setminus \frac{\nu^2-1}{\nu}, \frac{\nu^2-2}{\nu})$. Setting $v = u_t - u_x$, with $(x, t) \in \overline{\mathcal{R}_\eta} \times (\overline{\mathcal{T}_\nu} \setminus \frac{\nu^2-1}{\nu}, \frac{\nu^2-2}{\nu})$, we have that;

$$v_t + v_x = 0$$

$$v^0 = u_t^0 - u_x^0, (***)$$

the first line of (***) holding when $(x, t) \in \overline{\mathcal{R}_\eta} \times (\overline{\mathcal{T}_\nu} \setminus \frac{\nu^2-1}{\nu}, \frac{\nu^2-2}{\nu})$. Given finite (x, t) , we define $H_{x,t} : \overline{\mathcal{T}_\xi} \rightarrow {}^*\mathcal{R}$ by ;

$$H_{x,t}(s) = v\left(x - \frac{[s\xi]}{\xi}, t - \frac{[s\xi]}{\xi}\right) = v\left(x - \frac{[s\xi]}{\xi}, \frac{[t\nu]}{\nu} - \frac{[s\xi]}{\xi}\right)$$

Using (***) and the following Lemmas, 0.7 and 0.8, for finite $\{x_0, t_0, s\}$, with $x(s) = -\frac{[s\xi]}{\xi}$ and $y(s) = -\frac{[s\xi]}{\xi}$, so that $x^D(s) = y^D(s) = -1$, we have that;

$$\frac{dH^s}{dt}_{x_0, t_0} \simeq \left(-\frac{\partial v}{\partial x} - \frac{\partial v}{\partial t}\right)\Big|_{x_0-s, t_0-s} \simeq 0$$

so that;

$$(u_t - u_x)(x_0, t_0) = v(x_0, t_0) = H_{x_0, t_0}(0) \simeq H_{x_0, t_0}(s_0)$$

(where s_0 is chosen so that $\frac{[s_0\xi]}{\xi} = \frac{[t_1\nu]}{\nu} \simeq \frac{[t_0\nu]}{\nu}$; writing $\xi = \frac{\nu}{\kappa}$, with κ infinite, we have $\frac{[s_0\xi]}{\xi} = \frac{[t_0\nu]}{\nu}$, iff $\frac{\kappa[s_0\nu]}{\nu} = \frac{[t_0\nu]}{\nu}$ iff $[\frac{s_0\nu}{\kappa}] = \frac{[t_0\nu]}{\kappa}$. Replacing $[t_0\nu]$ with $[t_0\nu] + r$, where $0 \leq r < \kappa$, and $[t_1\nu] = [t_0\nu] + r$, so that $\kappa[t_1\nu]$ and, as $|\frac{r}{\nu}| < |\frac{\kappa}{\nu}| = |\frac{1}{\xi}| \simeq 0$, we must have $\frac{[t_0\nu]}{\nu} \simeq \frac{[t_1\nu]}{\nu}$. We can then solve $[\frac{s_0\nu}{\kappa}] = \frac{[t_1\nu]}{\kappa}$, taking $\frac{[t_1\nu]}{\nu} \leq s_0 < \frac{[t_1\nu]+\kappa}{\nu}$)

$$= v\left(x_0 - \frac{[s_0\xi]}{\xi}, \frac{[t_0\nu]}{\nu} - \frac{[s_0\xi]}{\xi}\right)$$

$$\simeq v\left(x_0 - \frac{[s_0\xi]}{\xi}, \frac{[t_1\nu]}{\nu} - \frac{[s_0\xi]}{\xi}\right)$$

$$\begin{aligned}
&= v(x_0 - \frac{[s_0\xi]}{\xi}, 0) \\
&= v(x_0 - \frac{[t_1\nu]}{\nu}, 0) \\
&\simeq v(x_0 - \frac{[t_0\nu]}{\nu}, 0)
\end{aligned}$$

Here, we have used the facts that v is S -continuous in t , and v^0 is S -continuous in x , consequent on $(u_t - u_x)$ S -continuous in t or u_{tt} and u_{tx} bounded, and $(u_t - u_x)^0$ is S -continuous in x or u_{tx}^0 and u_{xx}^0 bounded. These results follow as $u_{tt} = u_{xx}$, both $u_{tx} = u_{xt}$ and u_{xx} are bounded by the result of the previous lemma.

$$\begin{aligned}
&= u_t - u_x(x_0 - \frac{[t_0\nu]}{\nu}, 0) \\
&= u_t^0(x_0 - \frac{[t_0\nu]}{\nu}) - u_x^0(x_0 - \frac{[t_0\nu]}{\nu}), (***)
\end{aligned}$$

Rewriting (***) as $(u_t - u_x)(x', t') \simeq h(x', t')$

$$\text{with } h(x', t') = u_t^0(x' - \frac{[t'\nu]}{\nu}) - u_x^0(x' - \frac{[t'\nu]}{\nu})$$

define $\Theta_{x', t'} : \overline{\mathcal{T}}_\xi \rightarrow {}^*\mathcal{R}$ by ;

$$\Theta_{x', t'}(s) = u(x' + \frac{[s\xi]}{\xi}, t' - \frac{[s\xi]}{\xi})$$

Then, using (***) ,the definition of h , and Lemmas 0.7 and 0.8 again, this time with $x(s) = \frac{[s\xi]}{\xi}$, so that $x^D(s) = 1$;

$$\begin{aligned}
\frac{d\Theta^s}{dt}{}_{x', t'} &\simeq (\frac{\partial u}{\partial x} - \frac{\partial u}{\partial t})|_{x'+\frac{[s\xi]}{\xi}, t'-\frac{[s\xi]}{\xi}} \\
&\simeq (\frac{\partial u}{\partial x} - \frac{\partial u}{\partial t})|_{x'+\frac{[s\nu]}{\nu}, t'-\frac{[s\nu]}{\nu}}
\end{aligned}$$

(S -continuity of u_x in x and u_t in t , consequent on u_{xx} and u_{tt} bounded, follows from the previous lemma.)

$$\begin{aligned}
&\simeq -h(x' + \frac{[s\nu]}{\nu}, t' - \frac{[s\nu]}{\nu}) \\
&= -u_t^0(x' + \frac{[s\nu]}{\nu} - \frac{[t'\nu] - [s\nu]}{\nu}) + u_x^0(x' + \frac{[s\nu]}{\nu} - \frac{[t'\nu] - [s\nu]}{\nu}) \\
&= -u_t^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu}) + u_x^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu}), (\dagger)
\end{aligned}$$

We have, as above, that;

$$\Theta_{x',t'}(0) = u(x', t')$$

$$\Theta_{x',t'}(t') \simeq u(x' + \frac{[t'\nu]}{\nu}, t' - \frac{[t'\nu]}{\nu}) = u(x' + \frac{[t'\nu]}{\nu}, 0), (** ** ** *)$$

Using Lemma 0.8, (** ** ** *) and (\dagger), we obtain;

$$\begin{aligned} & u(x' + \frac{[t'\nu]}{\nu}, 0) - u(x', t') \\ & \simeq \Theta_{x',t'}(t') - \Theta_{x',t'}(0) \\ & = \int_{[0, \frac{[t'\nu]}{\nu})} \frac{d\Theta_{x',t'}^w}{dt} d\mu_\xi(w) \\ & \simeq \int_{[0, \frac{[t'\nu]}{\nu})} (-u_t^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu}) + u_x^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu})) d\mu_\xi(s) \end{aligned}$$

so that, using Lemma 0.5 and Lemma 0.9;

$$\begin{aligned} & u(x', t') \\ & \simeq u(x' + \frac{[t'\nu]}{\nu}, 0) + \int_{[0, \frac{[t'\nu]}{\nu})} (u_t^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \\ & \quad - \int_{[0, \frac{[t'\nu]}{\nu})} u_x^0(x' + \frac{2[s\nu]}{\nu} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \\ & = u(x' + \frac{[t'\nu]}{\nu}, 0) + \int_{[0, \frac{[t'\nu]}{\nu})} (u_t^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \\ & \quad - \int_{[0, \frac{[t'\nu]}{\nu})} u_x^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \\ & = u(x' + \frac{[t'\nu]}{\nu}, 0) + \int_{[0, \frac{[t'\nu]}{\nu})} (u_t^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \\ & \quad - \int_{[0, \frac{[t'\nu]}{\nu})} u_x^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) \end{aligned}$$

Now let $s = h_\rho(u) = \frac{[\rho u] + [t'\nu] - x'}{2}$, where $\xi = \kappa\rho$ and $\kappa \in {}^*\mathcal{N}$ is infinite, then, letting $w_{1,\xi}(s) = u_t^0(x' + 2\frac{[s\xi]}{\xi} - \frac{[t'\nu]}{\nu})$, we have, by refinement, that;

$$(w_{1,\xi} \circ h_\rho)(u) = u_t^0(x' + 2(\frac{[\rho u] + [t'\nu] - x'}{2}) - \frac{[t'\nu]}{\nu}) = u_t^0(\frac{[\rho u]}{\rho})$$

$$h_\rho^D(u) = \frac{1}{2}$$

and, using Lemma 0.9;

$$\begin{aligned} \int_{[0, \frac{[t'\nu]}{\nu}]} (u_t^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu})) d\mu_\xi(s) &= \frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_t^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \\ - \int_{[0, \frac{[t'\nu]}{\nu}]} u_x^0(x' + \frac{2[s\xi]}{\xi} - \frac{[t'\nu]}{\nu}) d\mu_\xi(s) &= -\frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_x^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \end{aligned}$$

so that;

$$\begin{aligned} u(x', t') &\simeq u(x' + \frac{[t'\nu]}{\nu}, 0) + \frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_t^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \\ &- \frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_x^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u), \quad (***) \end{aligned}$$

Using Lemma 0.10, we have that;

$$\begin{aligned} &\frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_x^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \\ &\simeq \frac{1}{2} (u^0(x' + \frac{[t'\nu]}{\nu}) - u^0(x' - \frac{[t'\nu]}{\nu})), \quad (****) \end{aligned}$$

Combining (***) and (****), we obtain that;

$$\begin{aligned} u(x', t') &\simeq u(x' + \frac{[t'\nu]}{\nu}, 0) - \frac{1}{2} (u^0(x' + \frac{[t'\nu]}{\nu}) - u^0(x' - \frac{[t'\nu]}{\nu})) \\ &+ \frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_t^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \\ &= \frac{1}{2} (u^0(x' + \frac{[t'\nu]}{\nu}) - u^0(x' - \frac{[t'\nu]}{\nu})) \\ &+ \frac{1}{2} \int_{[x' - \frac{[t'\nu]}{\nu}, x' + \frac{[t'\nu]}{\nu}]} u_t^0(\frac{[\rho u]}{\rho}) d\mu_\rho(u) \\ &\simeq \frac{1}{2} (u^0(x' + t') - u^0(x' - t')) \\ &+ \frac{1}{2} \int_{[x' - t', x' + t']} u_t^0(w) d\mu_\eta(w) \end{aligned}$$

as required. □

Remarks 0.4. *For finite values, this result matches d'Alembert's solution from the standard case, so the nonstandard wave equation is a reasonable model for pursuing a diffusion approach. This would help to understand photon paths, for example, in Maxwell's equation for zero charge and density.*

Lemma 0.5. *Let $F : {}^*\mathcal{R} \rightarrow {}^*\mathcal{R}$ be internal, and $\eta = \kappa\xi$ with $\kappa \in {}^*\mathcal{N}$, then;*

$$((F)_\eta)_\xi = (F)_\xi$$

Proof. Let $x = \frac{i}{\xi}$, with $i \in {}^*\mathcal{Z}$, then;

$$\begin{aligned} ((F)_\eta)_\xi(x) &= F_\eta(x) = F_\eta\left(\frac{i}{\xi}\right) = F_\eta\left(\frac{i\kappa}{\eta}\right) \\ &= F\left(\frac{i\kappa}{\eta}\right) = F\left(\frac{i}{\xi}\right) = (F)_\xi\left(\frac{i}{\xi}\right) = (F)_\xi(x) \end{aligned}$$

As both $((F)_\eta)_\xi$ and $(F)_\xi$ are μ_ξ measurable, this proves the result. \square

Lemma 0.6. *Let $f \in V(\overline{\mathcal{R}_\eta})$, with f^D S -continuous and bounded, then if $\epsilon \simeq 0$, with $\epsilon\sqrt{\eta}$ infinite, we have that;*

$$f^{D_\epsilon} \simeq f^D$$

Proof. Let $\epsilon = \frac{n}{\sqrt{\eta}} + \delta$, with $0 \leq \delta < \frac{1}{\sqrt{\eta}}$, and $n \in {}^*\mathcal{N}$ infinite, then;

$$\begin{aligned} f^{D_\epsilon}(x) &= \frac{f(x+\epsilon) - f(x)}{\epsilon} \\ &= \frac{f(x + \frac{n}{\sqrt{\eta}}) - f(x)}{\frac{n}{\sqrt{\eta}} + \delta}, \text{ as } f \in V(\overline{\mathcal{R}_\eta}) \end{aligned}$$

Without loss of generality, we have that;

$$\frac{n}{n+1} \frac{f(x + \frac{n}{\sqrt{\eta}}) - f(x)}{\frac{n}{\sqrt{\eta}}} < \frac{f(x + \frac{n}{\sqrt{\eta}}) - f(x)}{\frac{n}{\sqrt{\eta}} + \delta} < \frac{f(x + \frac{n}{\sqrt{\eta}}) - f(x)}{\frac{n}{\sqrt{\eta}}}$$

and;

$$\begin{aligned} &\frac{f(x + \frac{n}{\sqrt{\eta}}) - f(x)}{\frac{n}{\sqrt{\eta}}} \\ &= \frac{1}{n} \sum_{k=0}^{n-1} \frac{f(x + \frac{k+1}{\sqrt{\eta}}) - f(x + \frac{k}{\sqrt{\eta}})}{\frac{1}{\sqrt{\eta}}} \\ &= \frac{1}{n} \sum_{k=0}^{n-1} f^D\left(x + \frac{k}{\sqrt{\eta}}\right) \\ &\simeq f^D(x) \end{aligned}$$

as f^D is S -continuous, so that;

$$f^D(x) \simeq \frac{n}{n+1} f^D(x) < f^{D_\epsilon}(x) < f^D(x)$$

as $\frac{n}{n+1} \simeq 1$ and $f^D(x)$ is finite, giving that $f^D(x) \simeq f^{D_\epsilon}(x)$

□

Lemma 0.7. *Let $f \in V(\overline{\mathcal{R}_\eta} \times \overline{\mathcal{T}_\nu})$, $x : \overline{\mathcal{T}_\xi} \rightarrow \overline{\mathcal{R}_\eta}$, $y : \overline{\mathcal{T}_\xi} \rightarrow \overline{\mathcal{T}_\nu}$, be measurable, with the forward derivatives $\{f_x^D, f_{xy}^D, f_t^D, x^D, y^D\}$ bounded. Then, if $H \in V(\overline{\mathcal{T}_\nu})$, with $H(s) = f(x(s), y(s))$, then;*

$$H^D(s) \simeq f_x^D(x(s))x^D(s) + f_y^D(y(s))y^D(s)$$

when D relative to ξ is taken so that $\frac{\sqrt{\eta}}{\xi}$ and $\frac{\nu}{\xi}$ are infinite.

Proof. We have that;

$$\begin{aligned} H^D(s) &= \nu(H(s + \frac{1}{\xi}) - H(s)) \\ &= \nu(f(x(s + \frac{1}{\xi}), y(s + \frac{1}{\xi})) - f(x(s), y(s))) \\ &= \nu(f(x(s + \frac{1}{\xi}), y(s + \frac{1}{\xi})) - f(x(s), y(s + \frac{1}{\xi}))) \\ &\quad + \nu(f(x(s), y(s + \frac{1}{\xi})) - f(x(s), y(s))) \\ &= \frac{(f(x(s + \frac{1}{\xi}), y(s + \frac{1}{\xi})) - f(x(s), y(s + \frac{1}{\xi})))}{x(s + \frac{1}{\xi}) - x(s)} \xi (x(s + \frac{1}{\xi}) - x(s)) \\ &\quad + \frac{(f(x(s), y(s + \frac{1}{\xi})) - f(x(s), y(s)))}{y(s + \frac{1}{\xi}) - y(s)} \xi (y(s + \frac{1}{\xi}) - y(s)) \\ &= \frac{(f(x(s + \frac{1}{\xi}), y(s + \frac{1}{\xi})) - f(x(s), y(s + \frac{1}{\xi})))}{x(s + \frac{1}{\xi}) - x(s)} x^D(s) + \frac{(f(x(s), y(s + \frac{1}{\xi})) - f(x(s), y(s)))}{y(s + \frac{1}{\xi}) - y(s)} y^D(s) \\ &= \frac{(f(x(s + \epsilon), y(s + \frac{1}{\xi})) - f(x(s), y(s + \frac{1}{\xi})))}{\epsilon} x^D(s) + \frac{(f(x(s), y(s + \delta)) - f(x(s), y(s)))}{\delta} y^D(s) \\ &= f_x^{D\epsilon}(x(s), y(s + \frac{1}{\xi}))x^D(s) + f_y^{D\delta}(x(s), y(s))y^D(s) \end{aligned}$$

with $\epsilon = x(s + \frac{1}{\xi}) - x(s) \simeq 0$ and $\delta = y(s + \frac{1}{\xi}) - y(s) \simeq 0$, as x^D and y^D are bounded.

We have $x^D(s) = \xi\epsilon = C$, $y^D(s) = \delta\epsilon = D$, with C, D finite, so that $\epsilon\sqrt{\eta} = \frac{C\sqrt{\eta}}{\xi}$ and $\delta\nu = \frac{D\nu}{\epsilon}$ are infinite, by the hypotheses. Applying the result of Lemma 0.6, we have $f_x^{D\epsilon} \simeq f_x^D$, and $f_y^{D\delta} \simeq f_y^D$, so;

$$H^D(s) \simeq f_x^D(x(s), y(s + \frac{1}{\xi}))x^D(s) + f_y^D(x(s), y(s))y^D(s)$$

As f_{xy}^D is bounded (finite position), f_x^D is S -continuous in y , so that;

$$H^D(s) \simeq f_x^D(x(s), y(s))x^D(s) + f_y^D(x(s), y(s))y^D(s)$$

as required. □

Lemma 0.8. *If $H \in V(\overline{\mathcal{T}}_\xi)$ has the property that;*

$$H^D(s) \simeq 0$$

for $s \in \overline{\mathcal{T}}_\xi$ with s finite, then;

$$H(0) \simeq H(s)$$

for $s \in \overline{\mathcal{T}}_\xi$ with s finite.

If $G, R \in V(\overline{\mathcal{T}}_\xi)$ have the property that;

$$G^D(s) \simeq R(s)$$

for $s \in \overline{\mathcal{T}}_\xi$ with s finite, then;

$$G(s) \simeq G(0) + \int_{[0, \frac{[s\xi]}{\xi})} R(w) d\mu_\xi(w)$$

for $s \in \overline{\mathcal{T}}_\xi$ with s finite.

Proof. Using the definition of D , we have that;

$$\begin{aligned} |H(s) - H(0)| &= \left| \frac{1}{\xi} \sum_{k=0}^{[s\xi]-1} \xi(H(\frac{k+1}{\xi}) - H(\frac{k}{\xi})) \right| \\ &= \left| \int_{[0, \frac{[s\xi]}{\xi})} H^D(w) d\mu_\xi(w) \right| \\ &\leq \int_{[0, \frac{[s\xi]}{\xi})} |H^D(w)| d\mu_\xi(w) \\ &\leq \epsilon \int_{[0, \frac{[s\xi]}{\xi})} d\mu_\xi(w) \\ &\leq \epsilon \frac{[s\xi]}{\xi} \end{aligned}$$

where $\epsilon \in \mathcal{R}_{>0}$ is arbitrary. As $\frac{[s\xi]}{\xi}$ is finite, we conclude that;

$$|H(s) - H(0)| \simeq 0$$

as required for the first result. For the second result;

$$\begin{aligned}
& |G(s) - G(0) - \int_{[0, \frac{[s\xi]}{\xi})} R(w) d\mu_\xi(w)| \\
&= |\frac{1}{\xi} * \sum_{k=0}^{[s\xi]-1} \xi (H(\frac{k+1}{\xi}) - H(\frac{k}{\xi})) - \int_{[0, \frac{[s\xi]}{\xi})} R(w) d\mu_\xi(w)| \\
&= |\int_{[0, \frac{[s\xi]}{\xi})} H^D(w) d\mu_\xi(w) - \int_{[0, \frac{[s\xi]}{\xi})} R(w) d\mu_\xi(w)| \\
&\leq \int_{[0, \frac{[s\xi]}{\xi})} |(H^D - R)(w)| d\mu_\xi(w) \\
&\leq \epsilon \int_{[0, \frac{[s\xi]}{\xi})} d\mu_\xi(w) \\
&\leq \epsilon \frac{[s\xi]}{\xi}
\end{aligned}$$

where $\epsilon \in \mathcal{R}_{>0}$ is arbitrary. As $\frac{[s\xi]}{\xi}$ is finite, we conclude that;

$$|G(s) - G(0) - \int_{[0, \frac{[s\xi]}{\xi})} R(w) d\mu_\xi(w)| \simeq 0 \text{ as required.}$$

□

Lemma 0.9. *Integration by Substitution* Suppose $\{a, b\} \subset \mathcal{R}$, and let $f \in C([a, b])$ be continuous with corresponding $f_\eta \in V(\overline{\mathcal{R}}_\eta \cap^* [a, b])$, suppose that $h : [h^{-1}(a), h^{-1}(b)] \rightarrow [a, b]$ is continuous and increasing, so invertible, with corresponding $h_\xi \in V(\overline{\mathcal{R}}_\xi \cap^* [h^{-1}(a), h^{-1}(b)])$, and forward derivative h_ξ^D , then;

$$\int_{\overline{\mathcal{R}}_\eta \cap^* [a, b]} f_\eta(x) d\mu_\eta(x) \simeq \int_{\overline{\mathcal{R}}_\xi \cap^* [h^{-1}(a), h^{-1}(b)]} (f_\eta \circ h_\xi)(y) h_\xi^D(y) d\mu_\xi(y)$$

provided $\frac{\eta}{\xi}$ is infinite.

Proof. For bounded $f \in V(\overline{\mathcal{R}}_\eta \cap^* [a, b])$, we define the measure $\mu_{\eta, l, f}$ by;

$$\mu_{\eta, l, f}([\frac{i}{\xi}, \frac{i+1}{\xi})) = (f_\eta \circ h_\xi)(\frac{i}{\xi}) \mu_\eta[\frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta}, \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta})$$

and extend linearly to *-finite unions of intervals in the corresponding *-algebra \mathfrak{B} . We have that $\mu_\eta[\frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta}, \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta})$

$$\begin{aligned}
&= \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta} - \frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta} \\
&= \left(\frac{h_\xi(\frac{i+1}{\xi}) - \delta}{\eta} \right) - \left(\frac{h_\xi(\frac{i}{\xi}) - \delta'}{\eta} \right)
\end{aligned}$$

$$= \left(\frac{h_\xi(\frac{i+1}{\xi}) - h_\xi(\frac{i}{\xi}) - \delta''}{\eta} \right)$$

where $0 \leq \delta < 1$, $0 \leq \delta' < 1$ and $0 \leq |\delta''| < 1$.

so that;

$$|\mu_\eta\left[\frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta}, \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta}\right] - (h_\xi(\frac{i+1}{\xi}) - h_\xi(\frac{i}{\xi}))| \leq \frac{|\delta''|}{\eta} \leq \frac{1}{\eta}$$

and;

$$\mu_\eta\left[\frac{[\eta h_\eta(\frac{i}{\eta})]}{\eta}, \frac{[\eta h_\eta(\frac{i+1}{\eta})]}{\eta}\right] \simeq (h_\eta(\frac{i+1}{\eta}) - h_\eta(\frac{i}{\eta}))$$

and, as f_η is bounded, with $|f_\eta| \leq C$;

$$\begin{aligned} & |f_\eta(h_\xi(\frac{i}{\xi}))\mu_\eta\left[\frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta}, \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta}\right] - f_\eta(h_\xi(\frac{i}{\xi}))(h_\xi(\frac{i+1}{\xi}) - h_\xi(\frac{i}{\xi}))| \\ & \leq \left| \frac{C\delta''}{\eta} \right| \\ & \leq \frac{C}{\eta} \simeq 0, (\dagger) \end{aligned}$$

In particular, letting $\kappa = [\xi g^{-1}(b)] - [\xi g^{-1}(a)]$, we have that $\frac{C\kappa}{\eta} \simeq 0$ if $\frac{\eta}{\xi}$ is infinite, ($\dagger\dagger$);

$$\begin{aligned} & \text{Then, using } (\dagger); \\ & f_\eta(h_\xi(\frac{i}{\xi}))\mu_\eta\left[\frac{[\eta h_\xi(\frac{i}{\xi})]}{\eta}, \frac{[\eta h_\xi(\frac{i+1}{\xi})]}{\eta}\right] \\ & \simeq f_\eta(h_\xi(\frac{i}{\xi}))(h_\xi(\frac{i+1}{\xi}) - h_\xi(\frac{i}{\xi})) \\ & = \frac{1}{\xi} f_\eta(h_\xi(\frac{i}{\xi})) h_{xi}^D(\frac{i}{\xi}), (*) \end{aligned}$$

We claim that for f as in the statement of the lemma;

$\int_{[a,b]} f_\eta d\mu_\eta = \int_{[h^{-1}(a), h^{-1}(b)]} d\mu_{\eta,l,f}$ (\dagger), in which case, we obtain the result, as, using ($*$) and ($\dagger\dagger$);

$$\begin{aligned} & \int_{[a,b]} f_\eta(x) d\mu_\eta(x) \\ & = \int_{[h^{-1}(a), h^{-1}(b)]} d\mu_{\eta,l,f} \end{aligned}$$

$$\begin{aligned}
&= {}^* \sum_{[\xi h^{-1}(a)] \leq i \leq [\xi h^{-1}(b)]} \mu_{\eta, l, f} \left(\left[\frac{i}{\xi}, \frac{i+1}{\xi} \right) \right) \\
&\simeq {}^* \sum_{[\xi h^{-1}(a)] \leq i \leq [\xi h^{-1}(b)]} f_{\eta} \left(h_{xi} \left(\frac{i}{\xi} \right) \right) \frac{1}{\xi} h_{\xi}^D \left(\frac{i}{\xi} \right) \\
&= \frac{1}{\xi} {}^* \sum_{[\xi h^{-1}(a)] \leq i \leq [\xi h^{-1}(b)]} f_{\eta} \left(h_{\xi} \left(\frac{i}{\xi} \right) \right) h_{\xi}^D \left(\frac{i}{\xi} \right) \\
&= \int_{[h^{-1}(a), h^{-1}(b))} f_{\eta} \left(h_{\xi}(y) \right) h_{\xi}^D(y) d\mu_{\xi}(y) \\
&= \int_{[h^{-1}(a), h^{-1}(b))} (f_{\eta} \circ h_{\xi})(y) h_{\xi}^D(y) d\mu_{\xi}(y)
\end{aligned}$$

In order to show (\dagger) , we first consider the case when $g = \chi_{(c,d)}$ with $(c, d) \subset [a, b]$. Then, using the argument above and the fact that $\chi_{(c,d)}$ is bounded;

$$\begin{aligned}
&\int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta, l, g_{\eta}} \\
&\simeq {}^* \sum_{[\eta h^{-1}(a)] \leq i \leq [\eta h^{-1}(b)]} (f \circ h_{\xi}) \left(\frac{i}{\xi} \right) (h_{\xi} \left(\frac{i+1}{\xi} \right) - h_{\xi} \left(\frac{i}{\xi} \right)) \\
&= {}^* \sum_{r_c \leq i \leq r_d} (h_{\xi} \left(\frac{i+1}{\xi} \right) - h_{\xi} \left(\frac{i}{\xi} \right)) \\
&= h_{\xi} \left(\frac{r_d+1}{\xi} \right) - h_{\xi} \left(\frac{r_c}{\xi} \right) \\
&\simeq (d - c) \\
&\simeq \int_{[a,b]} \chi_{(c,d)} d\mu_{\eta} \\
&= \int_{[a,b]} f d\mu_{\eta}.
\end{aligned}$$

as h is continuous, where $r_c = \min(\{i : h_{\xi}(\frac{i}{\xi}) \geq c\})$, $r_d = \max(\{i : h_{\xi}(\frac{i}{\xi}) \leq d\})$,

By a similar argument, we can show that, if $f_r = \lambda_1 \chi_{(c_1, d_1)} + \dots + \lambda_r \chi_{(c_r, d_r)}$ is a finite combination of characteristic functions, then;

$$\int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta, l, f_{r, \eta}} \simeq \int_{[a,b]} f_{r, \eta} d\mu_{\eta} \quad (\#)$$

Now, considering the case when f is continuous on $[a, b]$, we can find a sequence $\{f_r : r \in \mathcal{N}\}$, such that $\lim_{r \rightarrow \infty} f_r = f$ pointwise, and each f_r has the property $(\#)$.

We claim first that $\lim_{r \rightarrow \infty} \int_{[a,b]} f_{r, \eta} d\mu_{\eta} = \int_{[a,b]} f_{\eta} d\mu_{\eta}$, $(\#\#)$. This follows as;

$$\begin{aligned}
& \lim_{r \rightarrow \infty} \circ \left(\int_{[a,b)} f_{r,\eta} d\mu_\eta \right) \\
&= \lim_{r \rightarrow \infty} \int_{[a,b)} \circ f_{r,\eta} dL(\mu_\eta) \\
&= \lim_{r \rightarrow \infty} \int_{[a,b)} f_r d\mu \\
&= \int_{[a,b)} f d\mu \\
&= \int_{[a,b)} \circ f_\eta dL(\mu_\eta) \\
&= \circ \left(\int_{[a,b)} f_\eta d\mu_\eta \right)
\end{aligned}$$

using the definition of the Riemann integral, together with the facts that f_η and $\{f_{r,\eta} : r \in \mathcal{N}\}$, are S -continuous, S -integrable, and piecewise S -continuous, S -integrable respectively.

We claim, secondly, that;

$$\lim_{r \rightarrow \infty} \circ \left(\int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,f_{r,\eta}} \right) = \circ \left(\int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,f_\eta} \right), \quad (\#\#\#)$$

This follows, by observing that;

$$\begin{aligned}
& \left| \int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,f_{r,\eta}} - \int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,f_\eta} \right| \\
&= \left| \sum_{[\xi h^{-1}(a)] \leq i \leq [\xi h^{-1}(b)]} (\mu_{\eta,l,f_{r,\eta}} \left(\left[\frac{i}{\xi}, \frac{i+1}{\xi} \right) \right) - \mu_{\eta,l,f_\eta} \left(\left[\frac{i}{\xi}, \frac{i+1}{\xi} \right) \right)) \right| \\
&\leq \sum_{[\xi h^{-1}(a)] \leq i \leq [\xi h^{-1}(b)]} |f_{r,\eta} \circ h_\xi \left(\frac{i}{\xi} \right) - f_\eta \circ h_\xi \left(\frac{i}{\xi} \right)| \left| h_\xi \left(\frac{i+1}{\xi} \right) - h_\xi \left(\frac{i}{\xi} \right) \right| \\
&= \int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,|f_\eta - f_{r,\eta}|} \\
&\leq C \epsilon_r
\end{aligned}$$

when $|f_\eta - f_{r,\eta}| \leq \epsilon_r$, and $C \simeq b - a$, where $\lim_{r \rightarrow \infty} \epsilon_r = 0$ if we assume uniform convergence. Now define $U(s_\eta) = \int_{[a,b)} s_\eta d\mu_\eta$, and $T(s_\eta) = \int_{[h^{-1}(a), h^{-1}(b))} d\mu_{\eta,l,s_\eta}$. Then, we have, using $(\#)$, $(\#\#)$, $(\#\#\#)$, that;

$$\begin{aligned}
& |U(f_\eta) - T(f_\eta)| \\
&= |U(f_\eta) - U(f_{r,\eta}) + U(f_{r,\eta}) - T(f_{r,\eta}) + T(f_{r,\eta}) - T(f_\eta)| \\
&\leq |U(f_\eta) - U(f_{r,\eta})| + |U(f_{r,\eta}) - T(f_{r,\eta})| + |T(f_{r,\eta}) - T(f_\eta)|
\end{aligned}$$

$$\leq \frac{\epsilon}{3} + \frac{\epsilon}{3} + \frac{\epsilon}{3} = \epsilon$$

As ϵ was arbitrary, we obtain the result. \square

Lemma 0.10. *Let $\{u, u^D, u^{D^2}\} \subset V(\overline{\mathcal{T}_\nu})$ be bounded, and let $\nu = \kappa\xi$, with $\kappa \in {}^*\mathcal{N}$ and κ infinite. Let $(u^D)_\xi \in V(\overline{\mathcal{T}_\xi})$ be the measurable counterpart of u^D . Then, if $\{a, b\} \subset \overline{\mathcal{T}_\xi}$ are finite ;*

$$\int_{(a,b)} (u^D)_\xi(x) d\mu_\xi(x) \simeq u(b) - u(a)$$

Proof. For $0 \leq j \leq \kappa - 1$, let $(u^D)_{\xi,j}(x) = \nu(u(x + \frac{j+1}{\nu}) - u(x + \frac{j}{\nu}))$. We claim that for $0 \leq j \leq \kappa - 2$, $(u^D)_{\xi,j} \simeq (u^D)_{\xi,j+1}$, (\dagger). We have that;

$$\begin{aligned} & |(u^D)_{\xi,j+1} - (u^D)_{\xi,j}| \\ &= |\nu(u(x + \frac{j+2}{\nu}) - u(x + \frac{j+1}{\nu})) - \nu(u(x + \frac{j+1}{\nu}) - u(x + \frac{j}{\nu}))| \\ &= |\nu(u(x + \frac{j+2}{\nu}) - 2u(x + \frac{j+1}{\nu}) + u(x + \frac{j}{\nu}))| \\ &= \frac{1}{\nu} |\nu^2(u(x + \frac{j+2}{\nu}) - 2u(x + \frac{j+1}{\nu}) + u(x + \frac{j}{\nu}))| \\ &= \frac{1}{\nu} |(u^{D^2}(x + \frac{j}{\nu}))| \\ &\leq \frac{C}{\nu} \end{aligned}$$

with $C \in \mathcal{R}$. By the triangle inequality, we have that, for $0 \leq j_1 \leq j_2 \leq \kappa - 1$;

$$|(u^D)_{\xi,j_2} - (u^D)_{\xi,j_1}| \leq \frac{C\kappa}{\nu} = \frac{C}{\xi} \simeq 0$$

It follow that;

$$(u^D)_\xi \simeq \frac{1}{\kappa} * \sum_{j=0}^{\kappa-1} (u^D)_{\xi,j}$$

and;

$$\begin{aligned} & \int_{(a,b)} (u^D)_\xi(x) d\mu_\xi(x) \\ & \simeq \int_{(a,b)} \frac{1}{\kappa} * \sum_{0 \leq j \leq \kappa-1} (u^D)_{\xi,j} \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{\kappa\xi} * \sum_{[a\xi] \leq i \leq [b\xi]} * \sum_{0 \leq j \leq \kappa-1} \nu(u(\frac{i}{\xi} + \frac{j+1}{\nu}) - u(\frac{i}{\xi} + \frac{j}{\nu})) \\
&= \frac{\nu}{\kappa\xi} (u(\frac{[b\xi]+1}{\xi}) - u(\frac{[a\xi]}{\xi})) \\
&= u(\frac{[b\xi]+1}{\xi}) - u(\frac{[a\xi]}{\xi}) \\
&\simeq u(b) - u(a)
\end{aligned}$$

as u is S -continuous.

□

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